



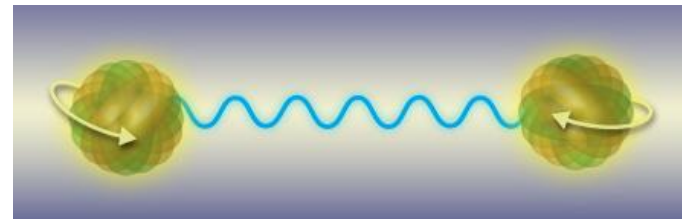
# Introduction – Quantum Optics and Entanglement

- **Quantum optics** refers to the study of **non-classical light** arising from quantized Maxwell's equations (single and few **photons**, **vacuum fluctuations**, **spontaneous emission**, etc.).
- Results in a fully **quantum-dynamical** model for both matter (e.g., electrons) and radiation (photons), which is necessary to study **quantum entanglement**.

# Introduction – Quantum Entanglement

- **Entanglement** is an experimentally verified property of nature where pairs of quantum systems are “connected” in some manner such that the quantum state of each system cannot be described independently.

$$|\psi\rangle = \frac{1}{\sqrt{2}} (|\uparrow\rangle_1 |\rightarrow\rangle_2 - |\rightarrow\rangle_1 |\uparrow\rangle_2)$$



<http://www.research.att.com>

- Measurements on one system of a pair of entangled systems collapses the wavefunction of the entangled system, so that the other system appears to “know” what measurement was performed on the first system, **instantaneously**.

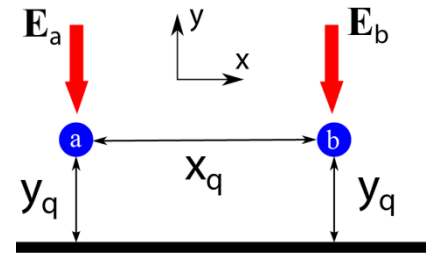
# Introduction – Quantum Entanglement

However, this does **not allow faster-than-light communications** (i.e, measurer #1 can't control what is measured, resulting in the no-communication theorem and no-cloning theorem).

So, what is entanglement good for?

- **Entanglement** is the cornerstone of much of **quantum computation** and quantum information theory.
  - **Generating, preserving, and controlling** entanglement is necessary for many quantum computer implementations.

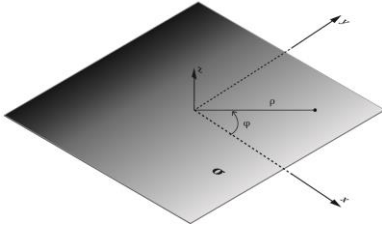
# Motivation



It is highly desirable to electronically manipulate the photonic spectrum of a **multi-level emitter** such as an **atom or quantum dot (QD)**, and to control entanglement, via a **macroscopic, easily-adjusted external parameter (e.g., bias)**.

**Surface plasmon polaritons (SPPs)** on graphene are **highly tunable**, and offer a promising way to achieve electronic control over a quantum emitter mediated by **graphene SPPs**.

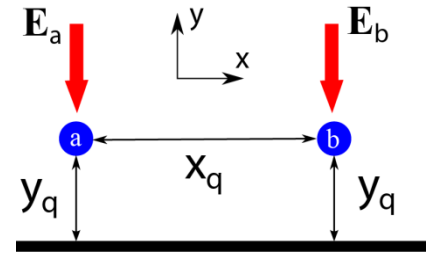
# Graphene – Electromagnetic Modeling



Infinite contiguous graphene sheet modeled as a two-sided impedance surface having conductivity  $\sigma$  (S).

$$\sigma(\omega) = \frac{ie^2k_B T}{\pi\hbar^2(\omega + i\Gamma_1)} \left( \frac{\mu_c}{k_B T} + 2 \ln \left( e^{-\frac{\mu_c}{k_B T}} + 1 \right) \right) + \frac{ie^2(\omega + i\Gamma_2)}{\pi\hbar^2} \int_0^\infty \frac{f_d(-\varepsilon) - f_d(\varepsilon)}{(\omega + i\Gamma_2)^2 - 4(\varepsilon/\hbar)^2} d\varepsilon$$

# Formulation



The **Hamiltonian** of the coupled system is the sum of QDs, pump, reservoir (graphene+vacuum), and their interaction

$$H = \int d\mathbf{r} \int_0^{+\infty} d\omega_\lambda \hbar\omega_\lambda \mathbf{b}^\dagger(\mathbf{r}, \omega_\lambda, t) \cdot \mathbf{b}(\mathbf{r}, \omega_\lambda, t) + \sum_{m=a,b} \hbar\omega_m \sigma_m^+(t) \sigma_m^-(t) - \sum_{m=a,b} (\sigma_m^+(t) + \sigma_m^-(t)) \mathbf{d} \cdot \mathbf{E}(\mathbf{r}_m, t)$$

$$\mathbf{E}(\mathbf{r}, t) = i \int_0^\infty d\omega_\lambda \int \sqrt{\frac{\text{Im}(\varepsilon(\mathbf{r}'; \omega_\lambda)) \hbar}{\pi \varepsilon_0}} \mathbf{G}(\mathbf{r}, \mathbf{r}'; \omega_\lambda) \cdot \mathbf{b}(\mathbf{r}', \omega_\lambda, t) d\mathbf{r}' + \text{H. C.}$$

$$\nabla \times \nabla \times \mathbf{G}(\mathbf{r}, \mathbf{r}') - k_0^2 \varepsilon(\mathbf{r}) \mathbf{G}(\mathbf{r}, \mathbf{r}') = k_0^2 \mathbf{I} \delta(\mathbf{r} - \mathbf{r}') \quad \text{Classical Green function}$$

$\sigma^+$  and  $\sigma^-$  are creation and annihilation operators for the atoms,  $\mathbf{b}$  are creation and annihilation (bosonic) operators for the photons.

# Formulation – Density Operator and Quantum Master Equation

Evolution of the system density matrix  $\rho_s = \sum_i p_i |\psi_i\rangle\langle\psi_i|$  is described by the Von Neumann equation,  $\partial_t \rho_s = -(i/\hbar)[H, \rho_s]$

$$\partial_t \rho_s(t) = -(i/\hbar)[\hat{V}, \rho_s(t)] + \mathcal{L}\rho_s(t), \quad \text{Evolution equation}$$

$$\hat{V} = -\hbar \sum_{j=a,b} (\Omega_j e^{-i\Delta_j t} \hat{\sigma}_j^\dagger + \Omega_j^* e^{i\Delta_j t} \hat{\sigma}_j) \quad \text{Source term}$$

$$\begin{aligned} \mathcal{L}\rho_s = & \sum_{i,j=a,b} \frac{\Gamma_{ij}(\omega_d)}{2} (2\hat{\sigma}_i \rho_s \hat{\sigma}_j^\dagger - \hat{\sigma}_i^\dagger \hat{\sigma}_j \rho_s - \rho_s \hat{\sigma}_i^\dagger \hat{\sigma}_j) \\ & + i[(g_{ab}(\omega_d) \hat{\sigma}_a^\dagger \hat{\sigma}_b + g_{ba}(\omega_d) \hat{\sigma}_b^\dagger \hat{\sigma}_a), \tilde{\rho}_s(t)], \end{aligned} \quad \mathcal{L}, \text{ Lindblad superoperator}$$

$\Omega_j = \mathbf{d} \cdot \mathbf{E}_j / \hbar$  (the effective Rabi frequency of the pump)



# Formulation – Density Operator and Quantum Master Equation

$$\Gamma_{ij}(\omega_d) = \frac{2\omega_d^2}{\varepsilon_0 \hbar c^2} \text{Im} \mathbf{d} \cdot \underline{\underline{\mathbf{G}}}(\mathbf{r}_i, \mathbf{r}_j, \omega_d) \cdot \mathbf{d},$$

$$g_{ij}(\omega_d) = \Delta_{ij} = \frac{\omega_d^2}{\varepsilon_0 \hbar c^2} \text{Re} \mathbf{d} \cdot \underline{\underline{\mathbf{G}}}(\mathbf{r}_i, \mathbf{r}_j, \omega_d) \cdot \mathbf{d},$$

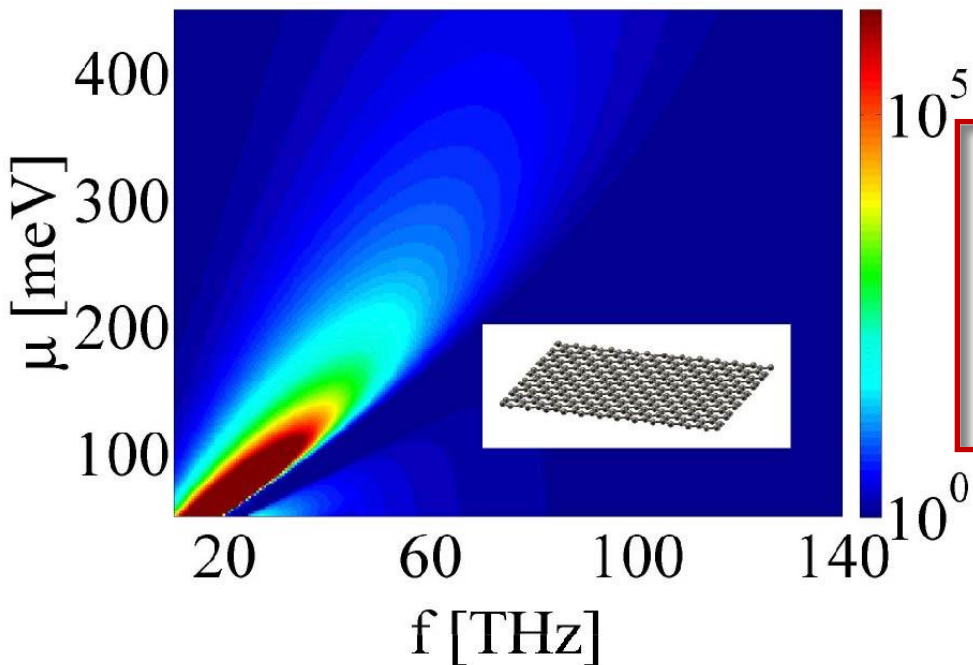
$\Gamma_{ii}$  is the rate of **spontaneous emission**, related to the **LDOS**

$\Gamma_{ij}$  ( $i \neq j$ ) is a **dissipative coupling term**

$g_{ij}$  is a **coherent dipole-dipole coupling parameter**

# Formulation – Computation of Green Functions

- In general, we use the commercial FDTD code **Lumerical** to numerically compute the **Green** function.
- Allows for a **true dipole source**.
  - Avoids numerical issues involved with discretizing a small line of current (e.g., CST).
  - Works very well for a variety of nanostructures (**plasmonic rods**, **grooves**, arrays of **nano-spheres**, **optical Yagi-Uda antennas**, etc.)
  - **Finite-sized graphene** problematic, so we assume infinite graphene.



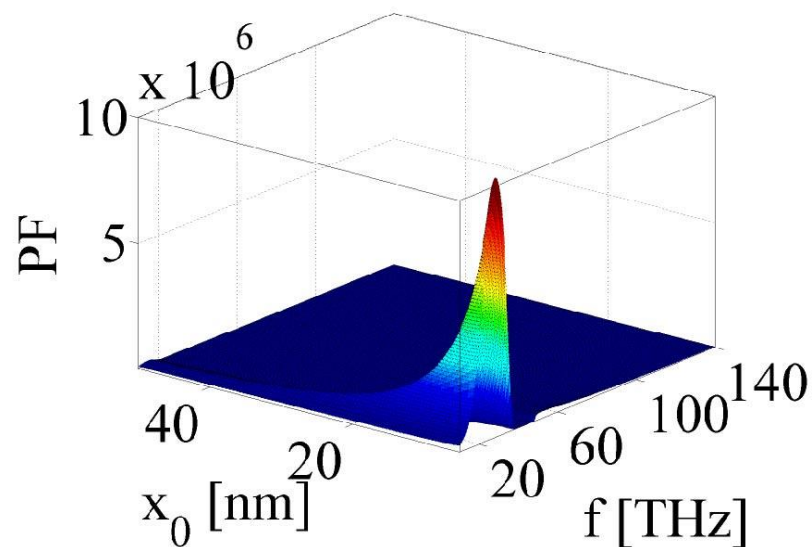
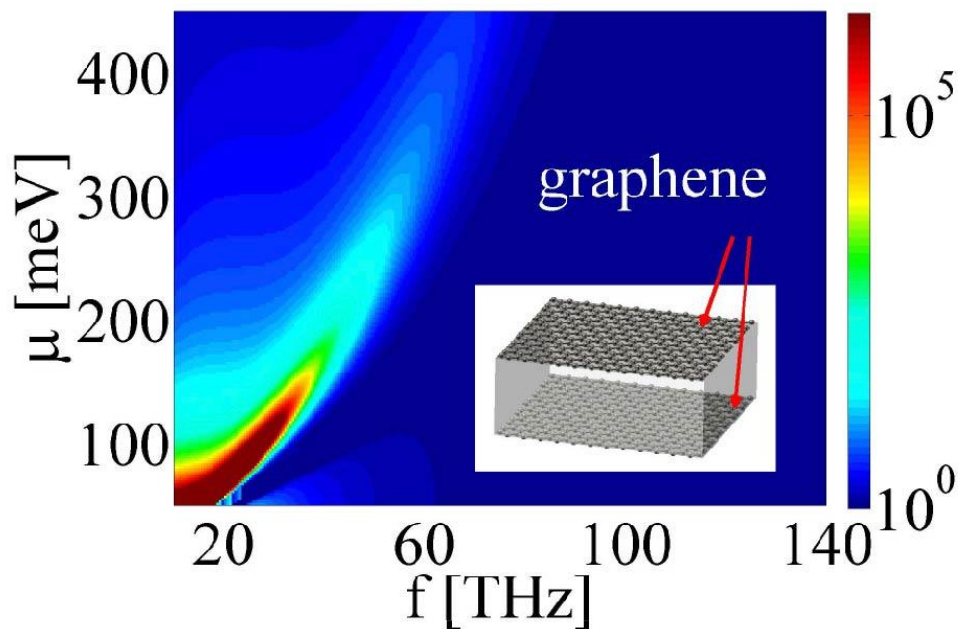
PHYSICAL REVIEW B **90**, 085414 (2014)

**Graphene as a tunable THz reservoir for shaping the Mollow triplet of an artificial atom via plasmonic effects**

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 (Received 29 April 2014; revised manuscript received 30 June 2014; published 12 August 2014)

**Purcell factor (PF),  $\Gamma_{ii}/\Gamma_0$**



# Entanglement

To assess entanglement we use the **concurrence**.

C=0: no entanglement,

C=1: maximum entanglement

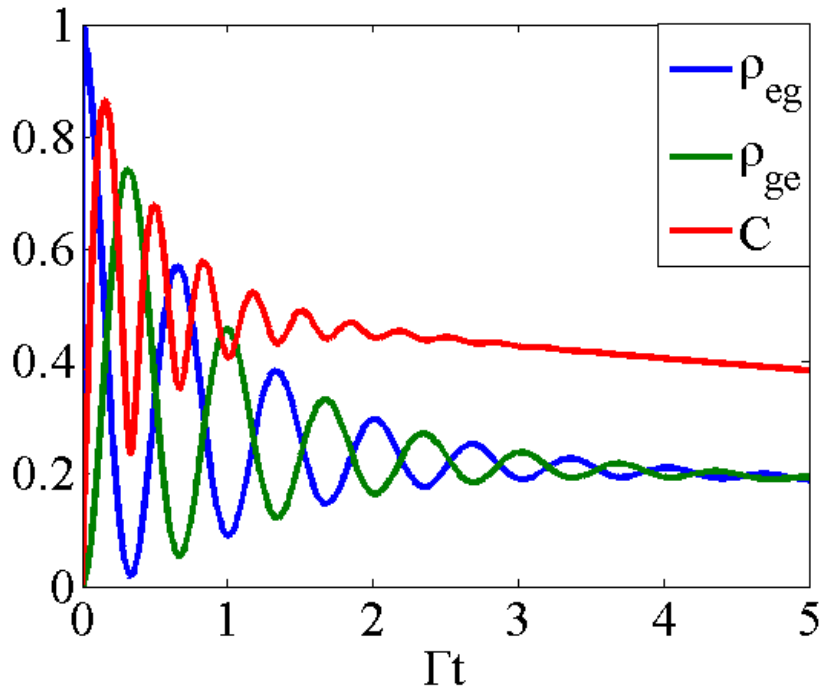
$$C = \max(0, \sqrt{u_1} - \sqrt{u_2} - \sqrt{u_3} - \sqrt{u_4})$$

where  $u_i$  are arranged in the descending order of the eigenvalues of the matrix  $\rho_s \tilde{\rho}_s$ ,

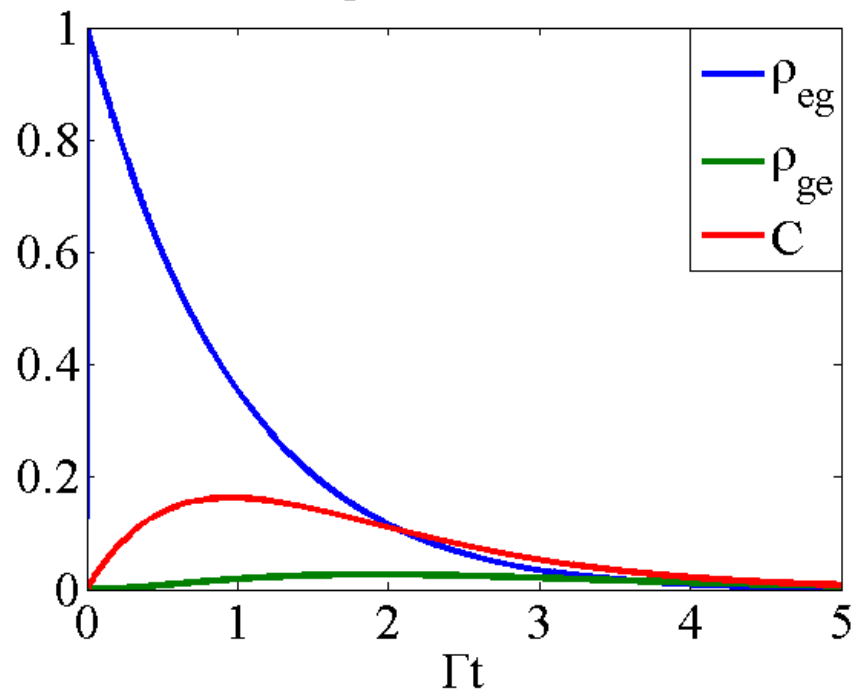
where  $\tilde{\rho}_s = \sigma_y \otimes \sigma_y \hat{\rho}_s^* \sigma_y \otimes \sigma_y$ ,  $\sigma_y$  is the Pauli matrix.

# Transient Entanglement via Spontaneous Emission in Vacuum

Separation =  $\lambda/12$



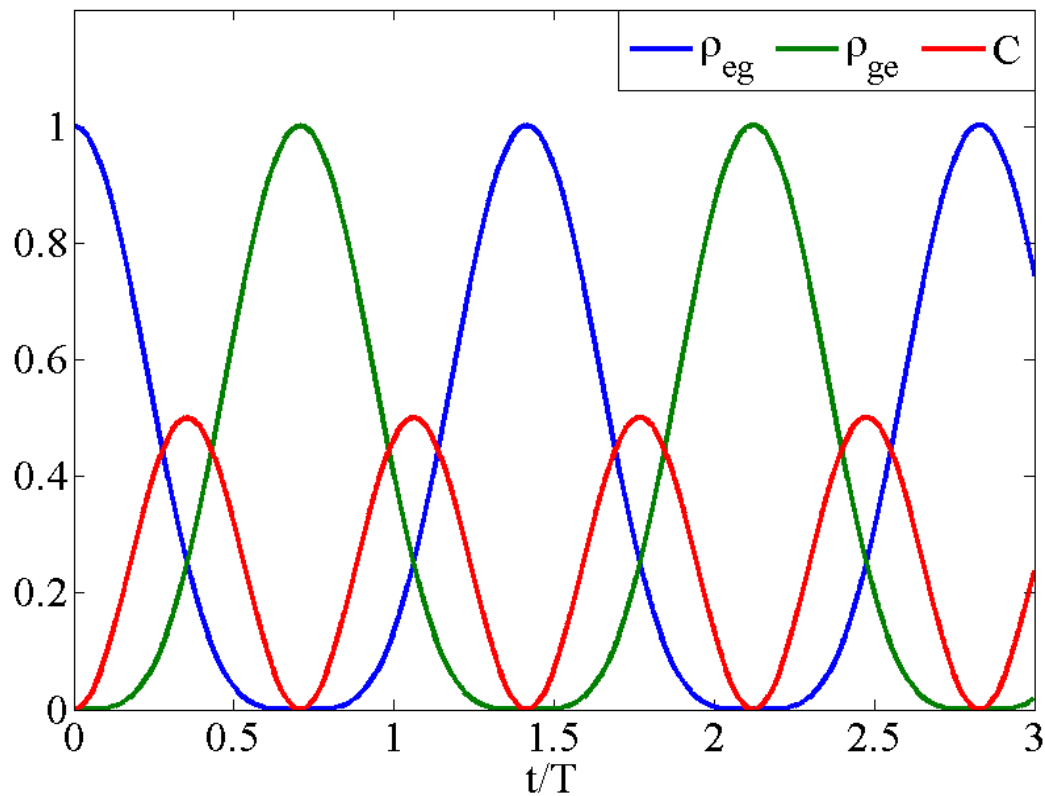
Separation =  $\lambda/2$



Population dynamics and **transient entanglement** between two quantum emitters in vacuum.

# Transient Entanglement via Spontaneous Emission in PEC Cavity

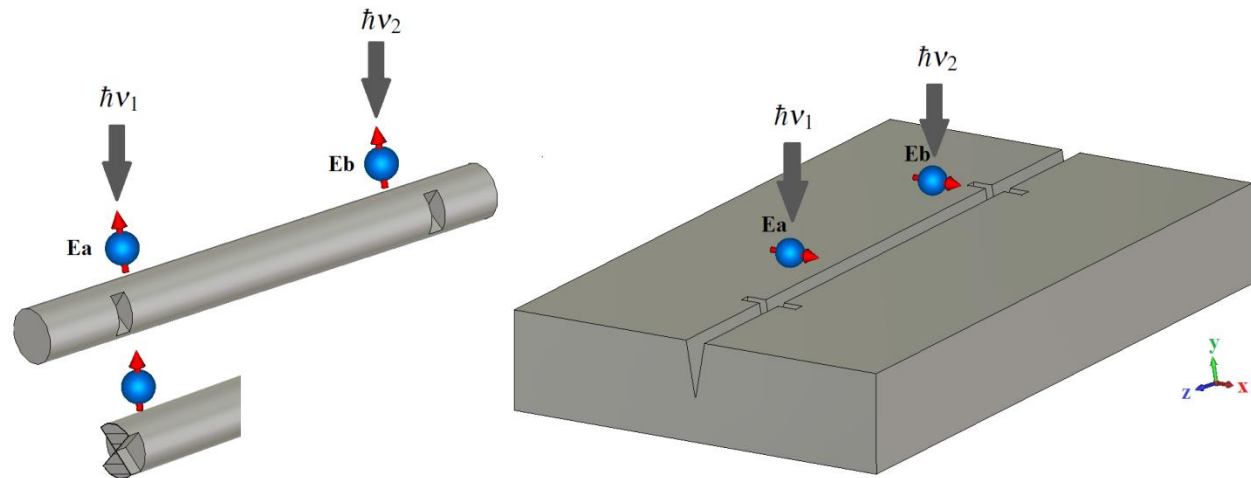
2 qubits inside an ideal PEC cavity



Population dynamics and transient entanglement between two quantum emitters in a PEC cavity, separation between emitters is  $\lambda/12$  at 80 THz.

# Graphene-Mediated Entanglement

- ◆ **Entanglement** is strong and relatively long-lived between closely-spaced emitters in vacuum.
- ◆ **Plasmonic** and other waveguiding systems can **aid entanglement** between far-separated emitters.

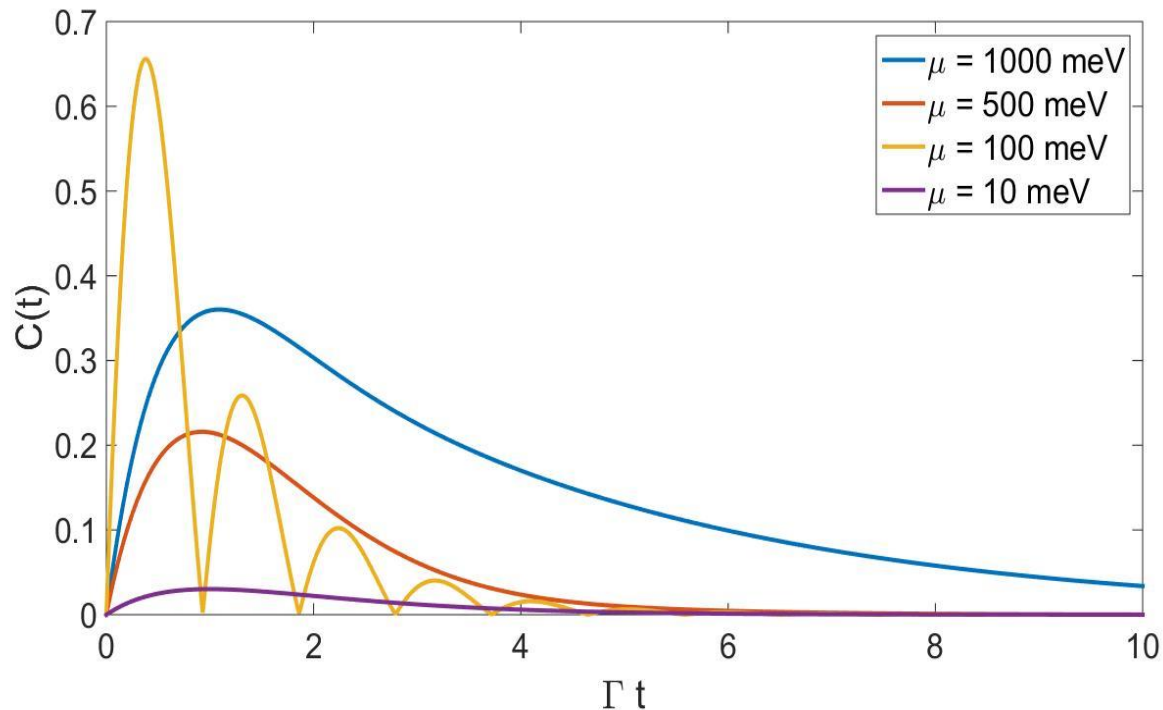


# Graphene-Mediated Entanglement

- ◆ However, we have found that graphene is **not useful for long-distance entanglement**
  - The graphene SPP is very tightly-confined to the surface, and, as a result,  $\lambda_{\text{SPP}}$  is **too small** ( $\lambda_{\text{SPP}} \sim \lambda_0/10$  to  $\lambda_0/100$ ) **for long-distance propagation.**
- ◆ Graphene does seem useful for **control** of entanglement.

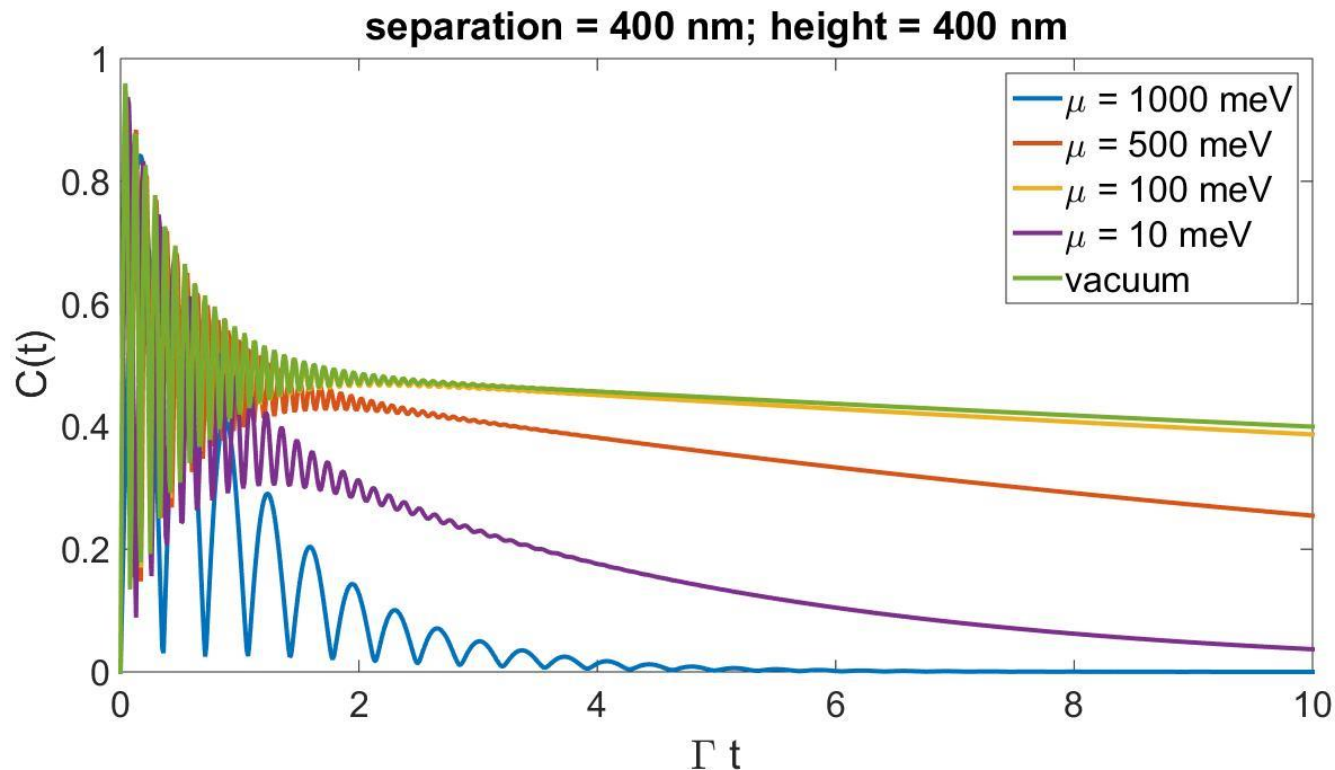


# Transient Entanglement via Spontaneous Emission over Graphene



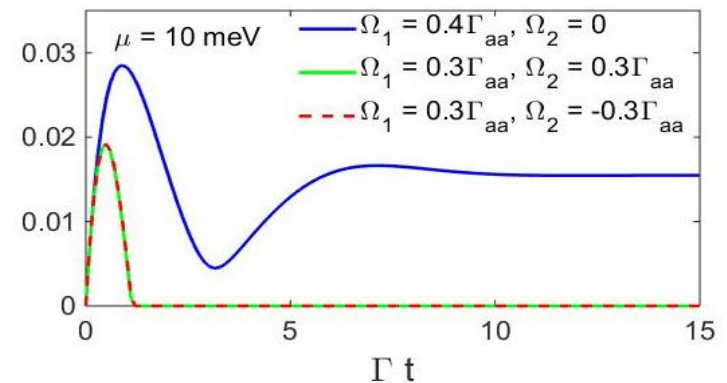
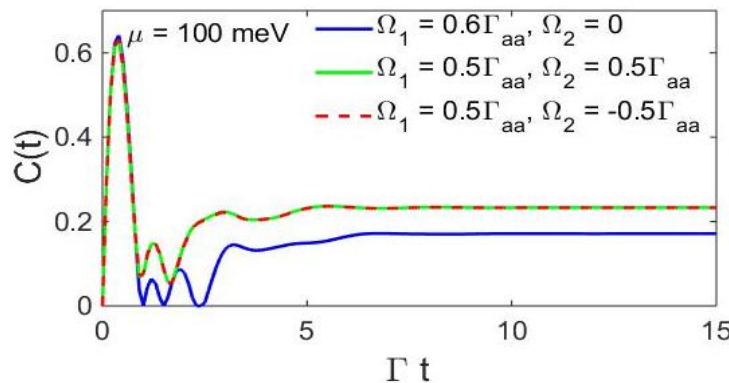
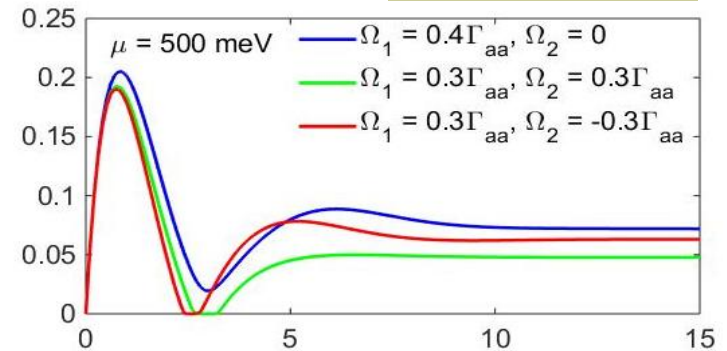
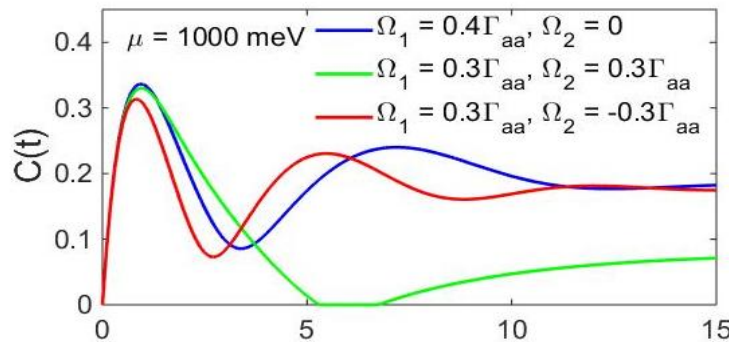
Transient entanglement between two quantum emitters placed at a distance 20 nm above graphene layer. Separation between emitters is equal to 100 nm. Frequency of the emitter dipole transition 40 THz.

# Control Via Graphene Bias of Long-Lived Transient Entanglement



Transient entanglement between two quantum emitters placed at a distance 400 nm above a graphene layer. Separation between emitters is equal to 400 nm. Frequency of the emitter dipole transition is 40 THz.

# Steady State Entanglement via External Pumping



Entanglement between two quantum emitters placed above graphene layer, which are pumped by external electromagnetic fields of intensities

$$\Omega_j = \mathbf{d} \cdot \mathbf{E}_j / \hbar \text{ (the effective Rabi frequency of the pump)}$$

# Conclusions

- Graphene is a promising material for quantum applications.
- **Tunability** of the graphene conductivity, and subsequent affects on SPPs, is a principle motivation for various applications related to **entanglement of quantum systems**.



Thank  
You!

Τηανκ  
Ψου!